

# Crystallographic Studies of the Biuret Reaction. I. Potassium *bis*-Biureto Cuprate(II) Tetrahydrate, $K_2[Cu(NHCONHCONH)_2] \cdot 4H_2O^*$

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The crystal structure of potassium *bis*-biureto cuprate(II) tetrahydrate,  $K_2[Cu(NHCONHCONH)_2] \cdot 4H_2O$ , has been determined, using the 3-dimensional X-ray diffraction intensities obtained with  $CuK\alpha$  radiation. Hydrogen atoms were located, and individual anisotropic thermal parameters were included in the refinement.

The copper atom in this complex is square-coordinated by four amide nitrogen atoms belonging to two (biuret)<sup>-</sup>-residues. The usual elongated octahedron about the copper atom is completed by two more distant nitrogen atoms, each of which belongs to one of the biuret residues in one of two neighbouring complexes.

The Cu-N-C angles at the two independent kinds of ligand nitrogen atoms are  $131^\circ$ . The configurations about these atoms seem to be distortedly trigonal. The complexes are not entirely planar, though the known dimensions of the biuret molecule are not altered significantly by coordination. The network of hydrogen-bonds is not particularly efficient.

The final *R* factor is 0.079 (including the unobserved reflections) and the standard deviations of the coordinates of the light atoms are about 0.007 Å.

## 1. Introduction

This investigation arose partly out of a general interest in metal-peptide inter-action (Freeman, Smith & Taylor, 1959) and partly out of a specific interest in the 'biuret reaction'.

It has been known for a long time that, when cupric ion is added to an alkaline solution of a protein or polypeptide, a violet-red colour is produced. This colour is reversibly changed to blue upon acidification of the solution. The reaction is a general one, but one of the simplest substances which undergo it is biuret,  $NH_2CONHCONH_2$ , (Brücke, 1883; Schiff, 1896) and the reaction has become identified with the name of this compound.

It was realized early that the violet-red colour was associated with the formation of a metal-peptide complex. Electrolysis of a copper-biuret solution showed the complex to be anionic (Ley & Werner, 1913; Rising, Hicks & Moerke, 1930), and Kober & Haw (1916; cf. Plekhan, 1952) ascribed the violet-red colour ( $\lambda_{max.} \approx 505 m\mu$ ) to the presence of four nitrogen atoms about the copper. No structural studies on the nature of the biuret-reaction complexes have, however, been reported.

## 2. Data

Potassium *bis*-biureto cuprate(II) tetrahydrate was prepared by shaking a fresh concentrated solution of cupric acetate (0.02 mole) with biuret (0.04 mole) dissolved in the calculated excess of aqueous potassium hydroxide (0.1 mole in 80 ml.). Violet-red crystals were obtained in 1-2 days after the addition of much alcohol (2 l.) containing a little alkali (Schiff, 1896, 1898). The crystals were unstable in air, but were adequately protected during the X-ray exposures by a thin film of collodion.

The space group was identified from systematic absences of reflections on Weissenberg photographs. The lattice constants were determined by fitting them to the spacings of six planes which gave reflections at large angles on Weissenberg photographs. These were taken with  $Cu K\alpha$  radiation on a large-radius camera and were calibrated by the superposition of the powder lines from a platinum wire on each film. An independent and identical value for  $d_{100}$  was obtained from a rotation photograph using the camera-adaptor and extrapolation procedure of Mathieson (1957).

$C_4H_{14}O_8N_6CuK_2$ ,  $M = 376.8$  ( $Cu K\alpha$ ,  $\lambda = 1.5418 \text{ \AA}$ ).

Monoclinic,

$$a = 3.843 \pm 0.004, \quad b = 13.321 \pm 0.008, \\ c = 13.916 \pm 0.008 \text{ \AA}; \quad \beta = 94^\circ 5' \pm 15'.$$

$U = 710.5 \text{ \AA}^3$ . For  $Z = 2$ ,  $D_c = 1.761 \text{ g.cm.}^{-3}$ .  
 $D_m$  not observed.  $F(000) = 414$ .

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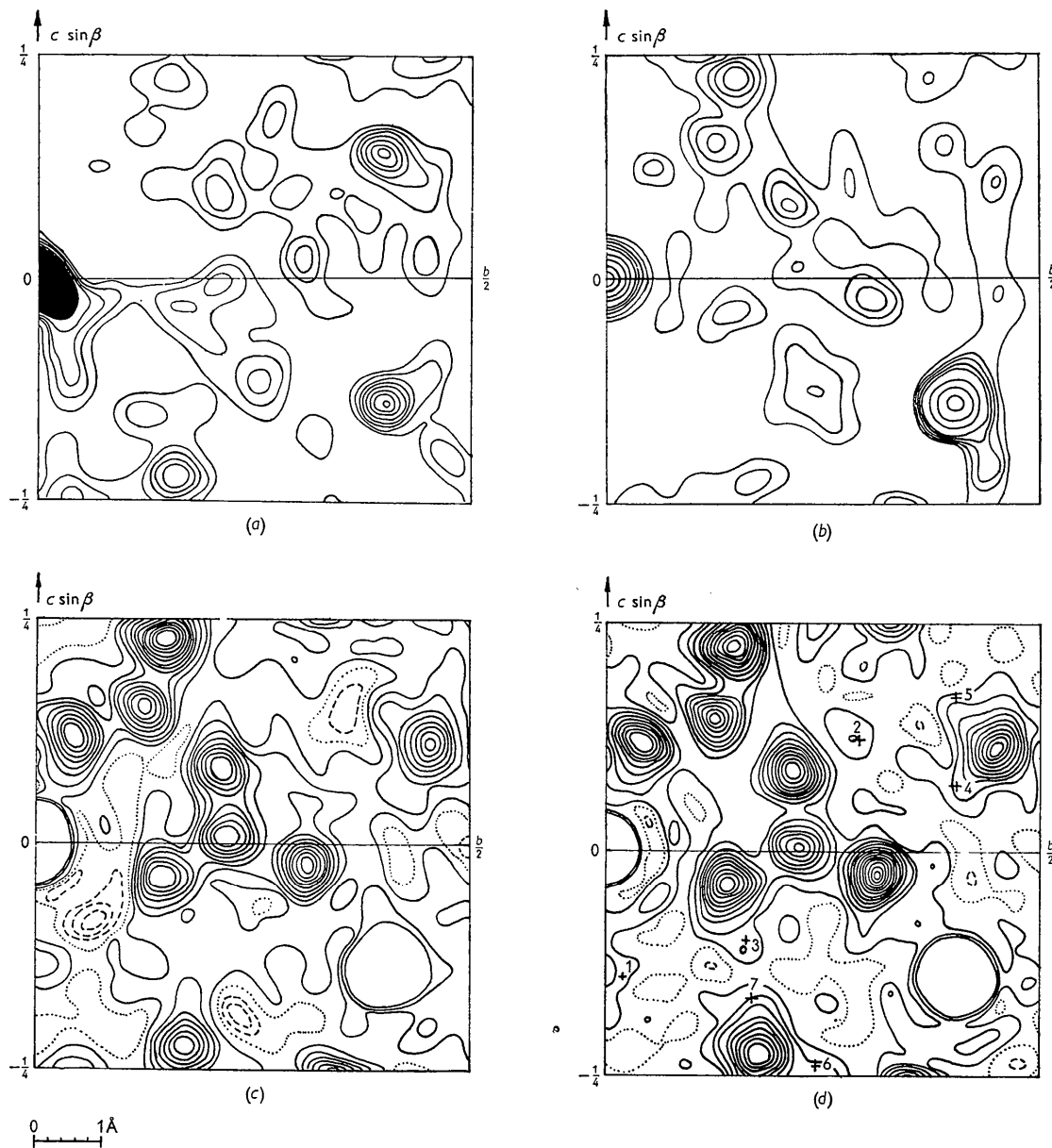


Fig. 1. Development of  $a$ -axis projection of  $\text{K}_2\text{Cu}(\text{NHCONHCONH})_2 \cdot 4 \text{H}_2\text{O}$ . (a) Scale arbitrary. (b) Scale arbitrary. After fourth contour, only one contour in four is shown. (c), (d) Contours at  $1 \text{ e.}\text{\AA}^{-2}$ , negative contours dashed, zero contour dotted. Crosses indicate probable hydrogen positions.

#### Systematic absences:

$h0l$  absent for  $(h+l)$  odd,

$0k0$  absent for  $k$  odd.

Space-group:  $P2_1/n-C_{2h}^5$ .

The intensities of 929  $0kl$ ,  $1kl$  and  $hkL$  ( $L=0$  to 8) reflections were estimated visually on Weissenberg photographs. In addition, 331 reflections in the same range were recorded as unobservably weak, and 38 of the  $3k1$ ,  $3k2$ ,  $4k1$  and  $4k2$  reflections were missed through an error in the setting of the Weissenberg

camera oscillation. Corrections to the observations were applied for the Lorentz, polarization and Tunell factors. The linear absorption coefficient of the crystals for  $\text{Cu } K\alpha$  radiation was  $80 \text{ cm.}^{-1}$ . The  $0kl$  and  $1kl$  data did not require absorption corrections since the crystal ( $0.3 \times 0.04 \times 0.04 \text{ mm.}^3$ ) used for these observations was rotated about its needle axis. For the  $hk0 \dots hk8$  photographs, another crystal ( $0.3 \times 0.1 \times 0.1 \text{ mm.}^3$ ) was rotated about an axis perpendicular to its long axis. In these cases the transmission factor varied from 10 to 45% and absorption was thus a serious problem.

Absorption corrections for the *c*-axis data were calculated *via* the method of Joel, Vera & Garaycochea (1953). Contours of these corrections, calculated at 40 values of  $(\alpha, \theta)$  per quadrant in a reciprocal lattice layer—where  $\alpha$  is the Weissenberg angular coordinate and  $\theta$  is the Bragg angle—were plotted on templates. The correction for each reflection could then be read to an estimated accuracy of 3% (of the calculable value) by placing the templates over the Weissenberg films.

An approximate scale factor for the *Ok*l data was obtained by Wilson's (1942) method. The other data were brought to the same scale *via* groups of common reflections. From a comparison of the final scaled  $F^2$  values of reflections observed about both axes it was concluded that the relative standard deviations of the intensities were about 20% for low intensities, 10% for medium intensities ( $100 < F^2 < 1600$ ) and 5% for strong reflections.

### 3. Structure determination

#### *Approximate structure*

The structure was solved in the projection along the short *a*-axis. Since the copper atom had to lie in a position of two-fold symmetry, an *Ok*l Fourier projection was calculated with all signs positive (Fig. 1(a)). This showed the location of the potassium ion. Successive projections showed, in turn, the biuret molecule (Fig. 1(b)) and the two water molecules (Fig. 1(c)) in the asymmetric unit. One further projection indicated small coordinate shifts, and the next corresponded to complete phase-determination in this zone (Fig. 1(d)).

One of the prime objects of this work was to discover whether the copper was bound to the biuret molecule at the amide nitrogen or oxygen atoms. In the structure-factor calculations associated with the above electron-density projections, all the atoms of the biuret molecule were initially given equal weights as nitrogen atoms. In the electron-density maps, however, the atoms of the amide groups from their first appearance had peak heights which left little doubt that nitrogen atoms were bonded to the central copper atom. The *Ok*l structure-factors calculated for this configuration of the ligand molecule, using positional coordinates read from the last Fourier projection with a temperature-factor  $\exp(-2.2 \sin^2 \theta / \lambda^2)$ , had a reliability-factor (for the observed reflections)  $R = \sum |F_o - F_c| / \sum |F_o| = 0.122$ . If either amide group was inverted so that it was coordinated to the metal *via* an oxygen atom, the agreement between the  $F_o$ 's and  $F_c$ 's became worse by about 1%.

An *Ok*l- $(F_o - F_c)$ -synthesis now not only indicated the expected changes in some coordinates and in the weights or temperature-parameters of the heavier atoms, but also had peaks corresponding to seven hydrogen atoms (Fig. 2). These provided strong evidence that the structure contained  $K^+$  and  $[Cu(NHCONHCONH)_2]^-$  ions and  $H_2O$  molecules

instead of, e.g.,  $K^+$  and  $OH^-$  ions and  $[Cu(NH_2CONHCONH_2)]^{++}$  complexes. While the important, tentative conclusions regarding the nature of the complex ion and the identities of the coordinating atoms were thus reached at an early stage of the work, they were substantiated by everything that happened in the further refinement of the structure.

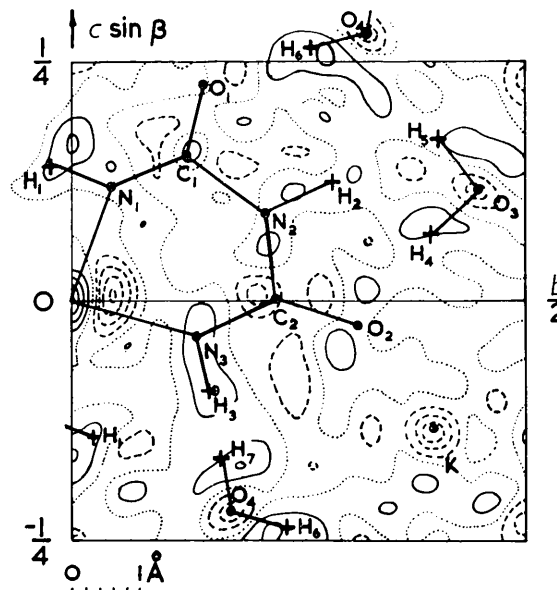


Fig. 2. *Ok*l( $F_o - F_c$ ) projection. Contours at  $0.5 \text{ e.}\text{\AA}^{-2}$ , negative contours dashed, zero contour dotted. Atomic positions from projection shown in Fig. 1(d), hydrogen atoms omitted from  $F_c$ .

The high intensity of the  $\bar{1}03$  reflection ( $|F_o| = 184$ ) made it likely that the chelating molecules lay in this plane. The angle between the  $\bar{1}03$  and  $(100)$  planes was  $38^\circ$ . The bond-lengths of the amide-groups in the *Ok*l-projection were compatible with those of normal amide groups tilted at this angle from the  $(100)$  plane. Approximate *x*-coordinates for the atoms of the biuret residues were calculated on this assumption. The potassium ion was placed at  $x = 0.5$  on the basis of a strong peak at  $(0.50, 0.80)$  on an *hk*0 Patterson projection, and of the reasonableness of the resultant  $K^+ \cdots O$  contacts with the oxygen atoms of the nearest biuret residues.

These *x*-coordinates were adequate for attaching signs to the 320 observed  $F(1kl)$ 's, whose scale at this stage was only approximate. Generalized projections  $C_1(yz)$  and  $S_1(yz)$  were computed from the first-layer data, and as a result the *x*-coordinates of the two  $H_2O$  molecules could also be approximately determined. When the  $H_2O$  molecules were included in the calculated  $F(1kl)$ 's,  $R$  fell from 0.33 to 0.29.

Further refinement of the combined zero- and first-layer *a*-axis data by generalized projections and least-squares quickly gave much improved agreement, the  $F(1kl)$ 's being scaled to their calculated values.

The  $x$ -coordinates of some atoms, however, converged on values which corresponded to non-planar amide-groups with unlikely bond-lengths. These anomalous structural features disappeared when comparison with the (by now) absorption-corrected  $c$ -axis data led to a 10% revision of the scale of the  $1kl$  intensities.

#### Refinement of coordinates

A least-squares refinement program for the Sydney University computer SILLIAC (Freeman, 1958) permitted the initial adjustments of the positional parameters to be carried out rapidly. While this program had provision for only one over-all two-parameter temperature factor  $\exp(-B \sin^2 \theta / \lambda^2 - \alpha h^2)$ , it had the virtues of availability and very high speed. The  $0kl$ ,  $1kl$  and  $hk0$  data were accordingly submitted to three cycles of refinement, which reduced  $R$  for these reflections from 0.147 to 0.113. The refinement process was now extended to the complete data. After five least-squares cycles, the over-all reliability factor  $R$  had dropped from 0.144 to 0.124. When the hydrogen atoms were included (see below),  $R$  was 0.120. The final temperature-parameters were  $B=2.24$  and  $\alpha=0.01$ . The greatest coordinate change in the last cycle was 0.003 Å in the  $x$ -direction, and the standard deviations of the positional parameters of the light atoms were between 0.02 and 0.03 Å.

The weighting system throughout this part of the refinement was that of Hughes (1941), i.e.

$$\begin{aligned} |w| &= 1/|F_o| \text{ for } |F_o| \geq 4|F_{\min.}|, \\ |w| &= \frac{1}{4}|F_{\min.}| \text{ for } |F_o| < 4|F_{\min.}| \text{ and for} \\ & \quad |F_c| > |F_{\text{unobs.}}|, \text{ and} \\ |w| &= 0 \text{ for } |F_c| \leq |F_{\text{unobs.}}|. \end{aligned}$$

( $|F_{\min.}|$  was the smallest observed  $F$ , and  $|F_{\text{unobs.}}|$  was the maximum value of an unobserved reflection.) With these weights,  $\sum w(\Delta F)^2$  was 50.5.

The parameters from the last SILLIAC least-squares cycle were supplemented by the calculation of individual isotropic temperature factors for all atoms except the hydrogens. For this purpose, Cochran's method (1951) was applied to an  $0kl$  ( $F_o - F_c$ )-projection. This step probably did not reduce the later computation. In view of the extent to which the

refinement of the anisotropic thermal parameters was ultimately pushed, a single isotropic temperature parameter would have been just as adequate a starting point. Seen in retrospect it would also have been a more correct early strategy, once the structure was established, to use all the available data instead of seeking economy by using only a limited number of  $F$ 's at the start.

The further refinement of the structure was carried out on a Datatron 205 computer. The program used only the diagonal terms of the least-squares matrix to evaluate the coordinate shifts, but for each atom a set of seven seven-term normal equations was solved to derive a scaling-factor and six coefficients for the anisotropic temperature factor

$$\exp(-4\beta'_i \sin^2 \theta) \exp \frac{1}{10}(\alpha'_i h^2 + \gamma'_i l^2 + \delta'_i hk + \epsilon'_i kl + \eta'_i hl).$$

The average of the scale factors for all atoms was substituted in all of the sets of normal equations which were then re-solved for the thermal parameters. [The final values of these parameters were converted to the parameters for the standard form of the temperature factor,

$$\exp(-\alpha_i h^2 + \beta_i k^2 + \gamma_i l^2 + \delta_i hk + \epsilon_i kl + \eta_i hl).$$

In this form they have been listed in Table 2.]

In this manner, seven refinement cycles were carried out. The weights attached to the  $F$ 's were the same as before, with the exception that all unobservably small reflections were given  $|w|=0$  for the first six cycles. The shifts indicated by an eighth cycle were not applied to the parameters; this cycle was used only to produce the final lists of structure-factors and the estimates of the various standard deviations.

The hydrogen atoms were included in the calculated structure-factors, but their coordinates and their common isotropic temperature factor were not automatically adjusted. After five cycles of refinement, the hydrogen positions were reviewed in the light of the changes which had occurred in the coordinates of the atoms to which they were presumed to be attached. This led to the discovery and correction of an arithmetical error in the original calculation of five of the hydrogen coordinates. As a result of this dis-

Table 1. *Final positional parameters, their standard deviations  $\sigma$ , and final thermal parameters for atomic temperature-factors*

$$T_i = \exp(-\alpha_i h^2 + \beta_i k^2 + \gamma_i l^2 + \delta_i hk + \epsilon_i kl + \eta_i hl)$$

Atom	$x_i$	$y_i$	$z_i$	$10^4 \sigma(x_i)$	$10^4 \sigma(y_i)$	$10^4 \sigma(z_i)$	$10^4 \alpha_i$	$10^4 \beta_i$	$10^4 \gamma_i$	$10^4 \delta_i$	$10^4 \epsilon_i$	$10^4 \eta_i$
C <sub>1</sub>	0.3882	0.1301	0.1469	17	5	6	384	27	21	26	-6	81
C <sub>2</sub>	0.0968	0.2226	0.0021	16	4	6	316	23	27	-37	11	69
N <sub>1</sub>	0.2800	0.0396	0.1146	16	4	5	535	23	31	-14	13	-80
N <sub>2</sub>	0.3034	0.2123	0.0893	15	4	5	467	21	32	12	1	-37
N <sub>3</sub>	-0.0251	0.1386	-0.0408	15	4	5	464	25	22	-9	-9	-3
O <sub>1</sub>	0.5557	0.1477	0.2251	14	4	4	591	33	19	-5	-25	-127
O <sub>2</sub>	0.0491	0.3125	-0.0272	14	4	4	685	23	20	26	5	-65
O <sub>3</sub>	0.9379	0.4469	0.1160	19	4	5	832	35	38	-32	5	56
O <sub>4</sub>	0.8621	0.3259	0.2786	17	5	5	716	50	43	57	-7	107
K <sup>+</sup>	0.5503	0.4030	-0.1387	4	1	1	417	35	26	18	3	-28
Cu	0.0000	0.0000	0.0000	—	—	—	458	20	18	12	-2	-52



Table 2 (cont.)

3,k,1		3,k,7		4,k,6		T,k,5		T,k,11		Z,k,2		Z,k,8		J,k,7			
0	262	274	0	51	-35	0	133	154	0	434	473	3	462	44	0	86	81
1	92	84	1	70	74	1	50	-32	1	149	135	4	429	421	1	91	102
2	132	119	2	128	123	2	49	72	2	249	274	5	49	48	2	194	195
3	356	348	3	203	217	3	57	-76	3	432	465	6	571	524	3	191	186
4	462	-38	4	180	192	4	223	210	4	536	532	7	61	61	4	275	301
5	282	424	5	209	214	5	442	-1	5	442	-1	8	183	170	5	112	-13
6	462	-38	6	209	214	6	176	171	6	134	119	9	264	-207	6	21	204
7	462	-38	7	454	-63	7	115	104	7	177	-186	10	64	58	7	80	71
8	457	28	8	101	101				8	278	283	11	181	149	8	142	127
9	462	31	9	118	-109				9	229	-215	12	182	166	9	120	122
10	465	-77	10	74	86				10	352	368	13	114	113	10	99	-90
11	465	-73	11	423	67				11	457	-18	14	148	126	11	112	19
12	462	-27							12	127	153	15	457	-4	12	62	84
13	457	22													13	93	98
3,k,2		3,k,8		4,k,7		T,k,6		T,k,12		Z,k,3		J,k,8		J,k,8			
1	325	333	1	226	208				1	230	244	1	601	659	1	182	197
2	179	183	2	58	-86				2	263	-277	2	306	319	2	166	172
3	168	173	3	77	24				3	169	157	3	415	462	3	161	120
4	116	97	4	49	-96				4	159	-129	4	453	-8	4	166	190
5	59	21	5	479	-25				5	315	360	5	169	-179	5	166	190
6	121	-115	6	479	-25				6	183	172	6	465	37	6	154	-169
7	92	113	7	477	12				7	495	38	7	470	40	7	143	112
8	66	-62	8	477	55				8	490	57	8	477	7	8	208	160
9	210	224	9	445	58				9	485	13	9	345	326	9	239	202
10	99	-100	10	219	208				10	485	46	10	345	164	10	176	182
11	140	132							11	99	92	11	252	214	11	232	173
12	116	137							12	107	-102	12	138	-107	12	108	-88
13	120	115							13	132	112	13	126	112	13	81	64
															13	67	-56
3,k,3		4,k,8		T,k,1		T,k,7		T,k,13		Z,k,4		J,k,3		T,k,3			
1	216	234	0	150	144				0	495	-117	0	400	365	1	252	253
2	316	-311	1	184	-179				1	185	188	1	188	183	2	64	96
3	195	214	2	149	163				2	400	84	2	298	293	3	464	26
4	177	-186	3	67	-76				3	157	75	3	157	75	4	463	29
5	209	232	4	49	45				4	347	350	4	161	191	5	62	-19
6	131	-130	5	67	-20				5	495	10	5	240	-249	6	117	-96
7	133	116	6	83	67				6	128	140	6	328	336	7	83	66
8	474	16	7	60	75				7	481	-92	7	452	-90	8	116	-104
9	77	69	8	87	80				8	461	82	8	405	383	9	149	115
10	152	158	9	104	101				9	485	-166	9	236	249	10	473	-63
11	452	-61	10	68	78				10	104	99	10	302	289	11	260	243
12	65	-68										11	220	211	12	119	-86
13	55	-89										12	465	36	13	294	220
3,k,4		4,k,1		T,k,2		T,k,8		T,k,14		Z,k,5		J,k,4		T,k,4			
1	65	-42	1	59	-53				1	202	232	1	113	127	0	104	103
2	200	196	2	83	55				2	495	13	2	298	-310	1	49	18
3	318	322	3	182	184				3	147	139	3	158	145	2	94	99
4	146	146	4	462	32				4	490	-82	4	490	-82	3	162	-148
5	265	275	5	171	191				5	481	62	5	158	145	4	59	-84
6	263	-245	6	171	191				6	481	62	6	158	145	5	78	74
7	211	204	7	176	170				7	110	-110	7	110	-110	6	129	128
8	465	-6	8	176	170				8	465	-31	8	465	-31	7	135	131
9	160	42	9	438	-19				9	231	251	9	231	251	8	129	128
10	451	-32													9	79	73
11	441	-13															
12	433	24															
3,k,5		4,k,2		T,k,3		T,k,9		T,k,15		Z,k,6		J,k,5		T,k,5			
0	407	427	0	256	260				0	490	4	0	485	497	1	113	127
1	54	58	1	450	44				1	484	-26	1	485	497	2	298	-310
2	161	161	2	180	180				2	485	48	2	158	145	3	158	145
3	176	148	3	142	144				3	485	48	3	490	-88	4	90	-88
4	123	-109	4	51	-69				4	485	48	4	159	147	5	127	112
5	201	-208	5	104	-108				5	485	48	5	159	147	6	127	112
6	63	83	6	104	-108				6	485	48	6	159	147	7	195	207
7	461	-77	7	104	-108				7	485	48	7	159	147	8	116	98
8	168	161	8	104	-108				8	485	48	8	159	147	9	64	-14
9	55	71	9	104	-108				9	485	48	9	159	147	10	117	90
10	315	280										10	166	151	11	113	111
11	436	-28										11	56	-62	12	104	-114
12	166	201										12	447	25	13	85	110
3,k,6		4,k,4		T,k,4		T,k,10		T,k,16		Z,k,7		J,k,6		T,k,6			
0	407	427	0	256	260				0	490	4	0	485	497	1	113	127
1	54	58	1	450	44				1	484	-26	1	485	497	2	298	-310
2	161	161	2	180	180				2	485	48	2	158	145	3	158	145
3	176	148	3	142	144				3	485	48	3	490	-88	4	90	-88
4	123	-109	4	51	-69				4	485	48	4	159	147	5	127	112
5	201	-208	5	104	-108				5	485	48	5	159	147	6	127	112
6	63	83	6	104	-108				6	485	48	6	159	147	7	195	207
7	461	-77	7	104	-108				7	485	48	7	159	147	8	116	98
8	168	161	8	104	-108				8	485	48	8	159	147	9	64	-14
9	55	71	9	104	-108				9	485	48	9	159	147	10	117	90
10	315	280										10	166	151	11	113	111
11	436	-28										11	56	-62	12	104	-114
12	166	201										12	447	25	13	85	110
3,k,6		4,k,4		T,k,4		T,k,10		T,k,16		Z,k,7		J,k,6		T,k,6			
1	111	100	0	458	-19				0	212	233	0	453	-77	0	453	-77
2	162	-150	1	458	-19				1	83	-99	1	338	334	1	338	334
3	176	166	2	458	-19				2	102	121	2	170	180	2	170	180
4	120	-115	3	458	-19				3	105	-68	3	174	187	3	174	187
5	102	104	4	458	-19				4	105	109	4	243	242	4	243	242
6	461	35	5	458	-19				5	481	941	5	458	41	5	458	41
7	306	288	6	458	-19				6	481	941	6	233	213	6	233	213
8	241	234	7	458	-19				7	481	941	7	463	-33	7	463	-33
9	452	49	8	458	-19				8	481	941	8	110	111	8	110	111
10	443	-22	9	458	-19				9	481	941	9	127	105	9	127	105
11	432	23										10	166	151	10	166	151
12	417	-71										11	56	-62	11	56	-62
3,k,6		4,k,4		T,k,4		T,k,10		T,k,16		Z,k,7		J,k,6		T,k,6			
1	1																

Table 3. *Positional parameters for hydrogen atoms*

Atom	$x_i$	$y_i$	$z_i$
H <sub>1</sub>	0.320	-0.010	0.170
H <sub>2</sub>	0.395	0.274	0.117
H <sub>3</sub>	-0.217	0.150	-0.094
H <sub>4</sub>	0.978	0.400	0.066
H <sub>5</sub>	0.911	0.405	0.172
H <sub>6</sub>	0.752	0.263	0.259
H <sub>7</sub>	0.733	0.337	0.334

Isotropic temperature-factor:  $\exp -(2.4 \sin^2 \theta / \lambda^2)$ .

would have been excessively optimistic to attach much significance to the actual peak positions. It was gratifying that the quality of the data was such that indications of the hydrogen atoms were found at all. Their persistence left little doubt that the points of attachment were correctly chosen—namely, H<sub>1</sub> attached to N<sub>1</sub>, H<sub>2</sub> to N<sub>2</sub>, H<sub>3</sub> to N<sub>3</sub>, H<sub>4</sub> and H<sub>5</sub> to O<sub>3</sub>, and H<sub>6</sub> and H<sub>7</sub> to O<sub>4</sub>. Reasonable coordinates could then be deduced.

From an analysis of the inter-atomic vectors it was unlikely that H<sub>1</sub>, H<sub>2</sub> and H<sub>3</sub> were involved in hydrogen bonds at all. H<sub>4</sub> lay on a short O<sub>3</sub>...O<sub>2</sub> vector (about 2.7 Å), H<sub>5</sub> on a short O<sub>3</sub>...O<sub>4</sub> vector (2.8 Å), H<sub>6</sub> between atoms of types O<sub>4</sub> and O<sub>1</sub> (2.8 Å), and H<sub>7</sub> between O<sub>4</sub> and N<sub>3</sub> (3.1 Å). There were no other short contacts except from N<sub>1</sub> to the water oxygen atom (type O<sub>4</sub>) nearest it (3.5 Å) and from N<sub>3</sub> to the O<sub>4</sub> atom of the adjacent unit cell (3.4 Å) in the  $\bar{x}$ -direction. If H<sub>1</sub> and H<sub>3</sub> were placed, respectively, on these two vectors then it turned out that they were also very close to the planes of the amide groups to which they belonged.

The hydrogen atoms were therefore assumed to lie on the above hydrogen bonds or lines of close approach, and to be at 1.0 Å from the oxygen atoms or 0.97 Å from the nitrogen atoms to which they were attached. This left only H<sub>2</sub>, which was placed in the plane of C<sub>1</sub>, N<sub>2</sub> and C<sub>2</sub> and roughly on the external bisector of C<sub>1</sub>N<sub>2</sub>C<sub>2</sub>. The coordinates of the hydrogen atoms were adjusted to maintain their relative positions with respect to the rest of the structure from cycle to cycle. The final coordinates are listed in Table 4.

#### Atomic scattering factors

For the light atoms, the curves of Berghuis *et al.* (1955) were used in the early calculations, and the Hoerni-Ibers (1954) atomic scattering factors in the final refinement. For potassium and copper, the K<sup>+</sup> and Cu<sup>+</sup> scattering curves of Berghuis *et al.* were corrected for anomalous dispersion by the method of Dauben & Templeton (1955). For copper, the size of this correction was about -2 electrons over the observable range, which was large compared with a slight uncertainty whether the copper should be treated as Cu<sup>0</sup>, Cu<sup>+</sup> or Cu<sup>++</sup>. The difference between these would be significant for only a few reflections at low  $\theta$ 's. Chemical intuition suggested that the charge originally on Cu<sup>++</sup> was partly neutralised by

chelation, the resultant charge of -2 being distributed over the peripheral atoms of the complex ion. The use of the Cu<sup>+</sup> curve was further justified by the absence of any peak at the origin of the first  $0kl(F_o - F_c)$ -projection (Fig. 2), and by the fact that in each cycle the 'atomic' scaling factor (see above) for copper was always close to the average value for the other atoms.

#### Rate of convergence

The root mean square changes of the parameters (of the atoms other than copper, potassium and hydrogen) during the eight final refinement cycles are shown in Table 4, together with other evidence of the progress of the calculation. For comparison we have also included the equivalent information for six cycles of least-squares refinement in which only the reflections with  $h+k+l=(2n+1)$ —i.e. those to which the copper atom made no contribution—were used. The purpose of this computational excursion was to see whether the refinement of such a structure could be effectively speeded up by carrying out a partial refinement with reflections minimally affected by a heavy atom whose position was already accurately known. (Note that the thermal parameter changes refer to the coefficients *before* conversion to the standard form.)

It is seen from Table 4 that convergence was achieved after *six* cycles of refinement with the *ca.* 400 *odd* reflections. While  $R$  and  $\Sigma w(\Delta F)^2$  were in fact smaller than the corresponding quantities for the same reflections treated as part of the complete data, the standard deviations of the atomic parameters were appreciably larger after the shortened refinement than those of the parameters obtained by the eight cycles using all 1260 reflections.

	$x_i$ (Å)	$y_i$ (Å)	$z_i$ (Å)
R.m.s. standard deviations of coordinates from:			
Complete data	0.006	0.007	0.006
Odd $F$ 's only	0.008	0.009	0.008
R.m.s. difference between final coordinates	0.009	0.010	0.013
Max. difference between final coordinates	0.018	0.021	0.022

Although the shortened refinement did not produce a value of any parameter which was significantly different from the 'complete data' value, some of the differences were considerably greater than the r.m.s. changes indicated by the final cycle of each sequence. While it is likely that the two sets of parameters would have been reconciled by following the shortened refinement with calculations using the complete data, we did not pursue the matter to this conclusion. It was apparent that this method offered little over-all economy in computations.

Table 4. Comparison of rates of convergence of refinements carried out, respectively, with the complete data (cycles 1-8) and with only the reflections with odd  $(h+k+l)$  (cycles 1'-6')

The r.m.s. values of the parameter-changes are calculated for light (C, N, O) atoms only

Refinement cycle	$10^4 \times$ r.m.s. value of			$10^3 \times$ r.m.s. value of						R%	$\Sigma w(\Delta F)^2$
	$\Delta x_i$	$\Delta y_i$	$\Delta z_i$	$\Delta \beta'_i$	$\Delta \alpha'_i$	$\Delta \gamma'_i$	$\Delta \delta'_i$	$\Delta \epsilon'_i$	$\Delta \eta'_i$		
1	51	8	8	47	199	9	17	11	10	12.3	50.54
2	29	5	6	22	84	6	14	17	5	9.5	24.06
3	20	3	4	18	35	5	20	12	4	9.0	19.30
4	14	2	3	10	51	2	17	15	2	8.7	18.34
5	6	1	2	9	20	2	7	4	1	8.6	17.53
6 <sup>a</sup>	14	3	4	9	30	1	21	7	3	7.9	17.23
7 <sup>b</sup>	10	2	2	10	21	2	10	5	3	8.1(9.4) <sup>d</sup>	15.20(7.37) <sup>d</sup>
8	4	1	1	4	15	2	5	2	1	7.9(9.1) <sup>d</sup>	14.45(6.90) <sup>d</sup>
1'	20	6	6	40	135	9	34	26	10	13.5	18.35
2'	13	3	5	14	25	5	33	18	5	11.3	13.49
3'	9	2	3	15	51	4	13	9	3	10.7	9.54
4'	5	2	1	8	31	3	11	9	3	9.9	8.22
5' <sup>c</sup>	8	2	2	5	38	2	9	5	2	9.3	7.32
6'	8	1	1	13	27	2	8	10	3	9.0	6.19

Superscripts: a. Error in H-coordinates detected and corrected. b. Unobserved  $F$ 's given non-zero weights. c. H-atoms included in  $F_c$ . d. Values in parentheses are for odd reflections alone

#### 4. Description of structure

We have used sub-script numerals to distinguish atoms of the same type within an asymmetric unit, and super-script numerals for atoms belonging to

asymmetric units other than that centred about the copper atom at (0, 0, 0). The relationships between the coordinates of such atoms and those listed in Table 1 are as follows:

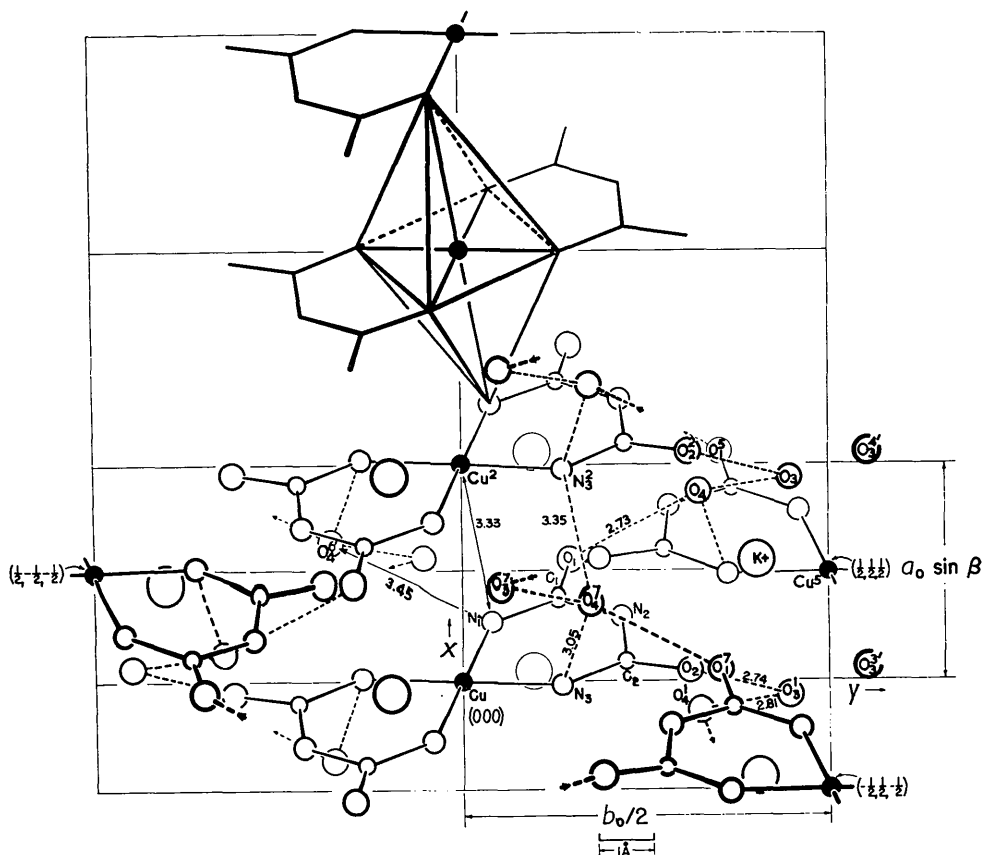


Fig. 3. Structure of  $K_2Cu(NHCONHCONH)_2 \cdot 4H_2O$  projected along c-axis.



- (2):  $x_i + 1, y_i, z_i$ ; (3):  $x_i + 1, y_i + 1, z_i$ ;  
 (4):  $x_i + 2, y_i + 1, z_i$ ; (5):  $x_i + \frac{1}{2}, \frac{1}{2} - y_i, z_i + \frac{1}{2}$ ;  
 (6):  $x_i + \frac{3}{2}, \frac{1}{2} - y_i, z_i + \frac{1}{2}$ ; (7):  $x_i - \frac{1}{2}, \frac{1}{2} - y_i, z_i - \frac{1}{2}$ ;  
 (8):  $x_i + \frac{1}{2}, \frac{1}{2} - y_i, z_i - \frac{1}{2}$ .

Dashes (') show atoms obtained from the corresponding un-dashed atoms by inversion through the copper atom of the same asymmetric unit. Thus  $O_3^5, O_4^5, O_3^{5'}$  and  $O_4^{5'}$  are the oxygen atoms of the four water molecules indicated by the molecular formula, considered with respect to the copper atom  $Cu^5$  at  $(\frac{1}{2}, \frac{1}{2}, \frac{1}{2})$ .

The structure is shown in projection along the  $z$ -axis in Fig. 3. Inter-atomic vectors and inter-vector angles are listed in Table 5. Safe estimates of the standard deviations are 0.015 Å for the bond-lengths and 2° for inter-bond angles.

Table 5. Inter-atomic distances and inter-vector angles

Bond-lengths		Other important distances	
Cu-N <sub>1</sub>	1.93 <sub>3</sub> Å	Cu-Cu <sup>2</sup>	3.84 Å
Cu-N <sub>3</sub>	1.93 <sub>2</sub>	N <sub>1</sub> -Cu <sup>2</sup>	3.33
C <sub>1</sub> -N <sub>1</sub>	1.34 <sub>2</sub>	N <sub>1</sub> -N <sub>3</sub> '	2.74
C <sub>1</sub> -N <sub>2</sub>	1.38 <sub>2</sub>	N <sub>1</sub> -N <sub>3</sub>	2.73
C <sub>1</sub> -O <sub>1</sub>	1.24 <sub>7</sub>	O <sub>4</sub> -N <sub>3</sub> <sup>6</sup>	3.35
C <sub>2</sub> -N <sub>2</sub>	1.40 <sub>9</sub>	O <sub>4</sub> -N <sub>3</sub> <sup>2</sup>	
C <sub>2</sub> -N <sub>3</sub>	1.33 <sub>7</sub>	O <sub>4</sub> -N <sub>1</sub> <sup>6'</sup>	3.45
C <sub>2</sub> -O <sub>2</sub>	1.27 <sub>4</sub>		
Hydrogen bonds		Environment of K <sup>+</sup> ion	
O <sub>1</sub> ...O <sub>4</sub>	2.73 Å	K <sup>+</sup> -O <sub>2</sub>	2.83 Å
O <sub>2</sub> <sup>2</sup> ...O <sub>3</sub>	2.74	K <sup>+</sup> -O <sub>2</sub> <sup>2</sup>	2.67
O <sub>3</sub> ...O <sub>4</sub>	2.81	K <sup>+</sup> -O <sub>1</sub> <sup>7</sup>	2.67
O <sub>4</sub> ...N <sub>3</sub> <sup>5</sup>	3.05	K <sup>+</sup> -O <sub>1</sub> <sup>8</sup>	2.89
O <sub>7</sub> <sup>4</sup> ...N <sub>3</sub>		K <sup>+</sup> -O <sub>3</sub> <sup>3'</sup>	2.78
		K <sup>+</sup> -O <sub>4</sub> <sup>4'</sup>	2.81
Angles between bonds, hydrogen-bonds and other vectors			
N <sub>1</sub> -Cu-N <sub>3</sub>	89.7°	C <sub>1</sub> -N <sub>2</sub> -C <sub>2</sub>	132.2°
N <sub>1</sub> -Cu <sup>2</sup> -N <sub>3</sub> <sup>2</sup>	87.9	C <sub>1</sub> -N <sub>2</sub> -H <sub>2</sub>	112.3
N <sub>1</sub> -Cu <sup>2</sup> -N <sub>1</sub> <sup>2</sup>	89.7	C <sub>2</sub> -N <sub>2</sub> -H <sub>2</sub>	115.4
Cu-N <sub>1</sub> -C <sub>1</sub>	131.5	N <sub>2</sub> -C <sub>2</sub> -O <sub>2</sub>	115.3
Cu-N <sub>1</sub> -Cu <sup>2</sup>	89.7	N <sub>2</sub> -C <sub>2</sub> -N <sub>3</sub>	117.5
C <sub>1</sub> -N <sub>1</sub> -Cu <sup>2</sup>	93.1	O <sub>2</sub> -C <sub>2</sub> -N <sub>3</sub>	127.3
Cu-N <sub>1</sub> -H <sub>1</sub>	119.5	C <sub>1</sub> -O <sub>1</sub> -O <sub>4</sub>	126.0
C <sub>1</sub> -N <sub>1</sub> -H <sub>1</sub>	107.6	C <sub>2</sub> <sup>2</sup> -O <sub>2</sub> <sup>2</sup> -O <sub>3</sub>	114.1
Cu-N <sub>3</sub> -C <sub>2</sub>	131.3	O <sub>2</sub> <sup>2</sup> -O <sub>3</sub> -O <sub>4</sub>	104.0
Cu-N <sub>3</sub> -O <sub>4</sub> <sup>7</sup>	111.9	O <sub>1</sub> -O <sub>4</sub> -N <sub>3</sub> <sup>5</sup>	97.7
C <sub>2</sub> -N <sub>3</sub> -O <sub>4</sub> <sup>7</sup>	93.6	O <sub>3</sub> -O <sub>4</sub> -N <sub>3</sub> <sup>5</sup>	131.3
Cu-N <sub>3</sub> -H <sub>3</sub>	112.0	O <sub>1</sub> -O <sub>4</sub> -N <sub>3</sub> <sup>6</sup>	126.4
Cu <sup>2</sup> -N <sub>3</sub> <sup>2</sup> -O <sub>4</sub> <sup>7</sup>	111.9	N <sub>3</sub> <sup>5</sup> -O <sub>4</sub> -N <sub>3</sub> <sup>6</sup>	73.6
C <sub>2</sub> -N <sub>3</sub> -H <sub>3</sub>	114.3	O <sub>3</sub> -O <sub>4</sub> -N <sub>3</sub> <sup>6</sup>	114.5
C <sub>2</sub> <sup>2</sup> -N <sub>3</sub> <sup>2</sup> -O <sub>4</sub> <sup>7</sup>	114.9	O <sub>1</sub> -O <sub>4</sub> -N <sub>1</sub> <sup>6'</sup>	170.3
N <sub>1</sub> -C <sub>1</sub> -N <sub>2</sub>	117.6	N <sub>3</sub> <sup>5</sup> -O <sub>4</sub> -N <sub>1</sub> <sup>6'</sup>	73.7
N <sub>1</sub> -C <sub>1</sub> -O <sub>1</sub>	126.1	O <sub>3</sub> -O <sub>4</sub> -N <sub>1</sub> <sup>6'</sup>	79.2
O <sub>1</sub> -C <sub>1</sub> -N <sub>2</sub>	116.3		

### Environment of copper atom

The copper atom is bonded to four nitrogen atoms lying at distances of 1.93 Å and at the corners of an almost perfect square of side 2.73 Å. Adjacent complexes are stacked along lines parallel to the  $x$ -axis, with their 'coordination-squares' at such an inclination that each bonded nitrogen atom N<sub>1</sub> also lies in the octahedral position with respect to the copper atom of one of the two adjacent complexes. The environment of the metal atom is thus comparable with that in copper phthalocyanine (Robertson, 1935, 1936, 1937). The octahedra of nitrogen atoms about Cu and Cu<sup>2</sup> share the N<sub>1</sub>-N<sub>1</sub><sup>2</sup> edge. The distance between the copper atom and the octahedral nitrogen is the minimum distance between any two points in adjacent complexes, but its size (3.33 Å) implies that any inter-action between these atoms is weak.

### Dimensions of ligand molecule

The biuret molecule has previously been studied in the crystals of biuret hydrate (Hughes, Yakel & Freeman, 1961) and *bis*-biuret cadmium(II) chloride (Cavalc, Nardelli & Fava, private communication). The dimensions found in the hydrate and in the copper complex are shown in Fig. 4; the following discussion applies substantially also to the cadmium derivative.

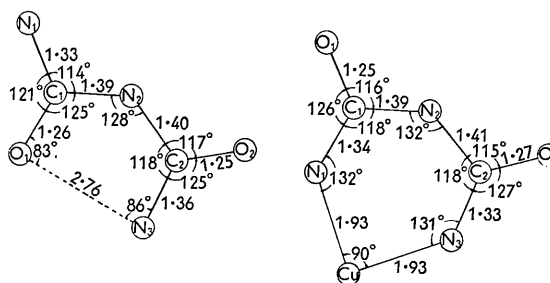


Fig. 4. Comparison of dimensions of biuret in (left) the hydrate and (right) the copper complex.

Coordination to copper seems to have no significant effect upon the linear and angular dimensions of the biuret molecule, although the configuration of the two amide groups is changed. There is thus no detectable change in the relative contributions of the resonance forms of the biuret molecule (Hughes *et al.*, 1961). Neither in biuret nor in the copper complex are there significant differences between the lengths of equivalent C-N bonds (C<sub>1</sub>-N<sub>1</sub> and C<sub>2</sub>-N<sub>3</sub>; C<sub>1</sub>-N<sub>2</sub> and C<sub>2</sub>-N<sub>2</sub>). In the complex, the difference between the bond-lengths of C<sub>1</sub>-O<sub>1</sub> and C<sub>2</sub>-O<sub>2</sub> lies just inside the limit of 'possible significance' according to the usual  $t_0$ -test (Cruickshank & Robertson, 1953).

The value of the angle C<sub>1</sub>-N<sub>2</sub>-C<sub>2</sub> is abnormally large in the complex (132°) as it is also in the free molecule (128°). The cause in both cases is almost certainly steric strain—arising in the free molecule from an internal hydrogen bond (Hughes *et al.*, 1961), and in the complex from the accommodation of the metal

Table 6. *Details of planes of best fit*

Each plane is represented by  $lx' + my' + nz' + k = 0$  with respect to orthogonal axes, relative to which the usual monoclinic coordinates of a point  $(x, y, z)$  become  $x' = x + z \cos \beta$ ,  $y' = y$ , and  $z' = z \sin \beta$

Plane	$l$	$m$	$n$	$k$	Angles between normal to plane and recip. axes			Angle between plane and plane 1
					$x^*$	$y^*$	$z^*$	
1	+0.8666	-0.1221	-0.4838	0	33.9°	97.0°	118.9°	—
2	+0.8885	-0.1123	-0.4450	+0.0896	—	—	—	2.7°
3	+0.8597	-0.0595	-0.5073	+0.1287	—	—	—	3.9
4	+0.8770	-0.0817	-0.4735	+0.0643	32.75	94.7	118.3	2.5

Atom	Deviations (in Å) from plane				Max. standard deviations of atomic positions (in Å)
	1	2	3	4	
Cu	0	—	—	—	—
C <sub>1</sub>	—	+0.0046	—	+0.0093	0.0075
O <sub>1</sub>	—	-0.0017	—	-0.0272	0.0055
N <sub>1</sub>	0	-0.0015	—	-0.0165	0.0066
N <sub>2</sub>	—	-0.0014	+0.0005	+0.0627	0.0066
C <sub>2</sub>	—	—	-0.0018	+0.0042	0.0075
O <sub>2</sub>	—	—	+0.0007	-0.0364	0.0055
N <sub>3</sub>	0	—	+0.0006	+0.0039	0.0066

Key to planes. Plane 1:  $\text{CuN}_1\text{N}_3(\text{N}'\text{N}'_3)$ . Plane 2:  $\text{C}_1\text{O}_1\text{N}_1\text{N}_2$ . Plane 3:  $\text{C}_2\text{O}_2\text{N}_2\text{N}_3$ . Plane 4:  $\text{C}_1\text{O}_1\text{N}_1\text{N}_2\text{C}_2\text{O}_2\text{N}_3$

atom. Similarly the internal Cu-N-C angles at both N<sub>1</sub> and N<sub>3</sub> are 131°, a considerable distortion from the value for trigonal bonds.

#### Configurations at ligand nitrogen atoms

Apart from its bonds with Cu and C<sub>2</sub>, the atom N<sub>3</sub> has two short contact distances, each with a water oxygen atom of type O<sub>4</sub> (see Fig. 3). These atoms complete a distorted tetrahedron about N<sub>3</sub>. We have assumed that the hydrogen-bond N<sub>3</sub> ··· H-O<sub>4</sub><sup>7</sup> (3.05 Å) involves the hydrogen atom H<sub>7</sub><sup>7</sup> of the water molecule. The amide hydrogen H<sub>3</sub> is then not used in any hydrogen bond: we have placed it on the second N<sub>3</sub> ··· O<sub>4</sub> vector (3.35 Å), which lies almost in the CuN<sub>3</sub>C<sub>2</sub> plane. Apart from its length, this vector is precluded from being a hydrogen bond by the angles which it makes with the three certain hydrogen-bonds at O<sub>4</sub>. It is equally possible that H<sub>3</sub> should have been placed along the bond N<sub>3</sub>-H ··· O<sub>4</sub><sup>7</sup>, and that H<sub>7</sub><sup>7</sup> lies roughly in the direction O<sub>4</sub><sup>7</sup> ··· N<sub>3</sub><sup>2</sup> (but at reasonable angles with the other bonds to O<sub>4</sub><sup>7</sup>). The positions of N<sub>3</sub> and O<sub>4</sub> would be scarcely affected, and this arrangement would leave N<sub>3</sub> with a more satisfying, distorted trigonal arrangement of three bonded contacts (bond-angles 131°, 112° and 94°).

The configuration about N<sub>1</sub> is similarly a distorted triangle, consisting in this case of Cu, C<sub>1</sub> and H<sub>1</sub>. The N<sub>1</sub>-H<sub>1</sub> bond does not take part in hydrogen bonding, and the most reasonable position for H<sub>1</sub> is close to the plane of the amide group. The copper atom Cu<sup>2</sup> of the adjacent complex lies almost exactly on the normal to both the Cu-N<sub>1</sub> and C<sub>1</sub>-N<sub>1</sub> bonds.

#### Planarity of amide groups

Just as the free biuret molecule is not planar, so the coordinately-bound biuret residue is bent slightly

about the line joining Cu to N<sub>2</sub>, the two amide groups then being rotated about their C-N<sub>2</sub> bonds until their normals are inclined at 4° 56' to each other. Neither of the two planar amide groups thus lies in the plane of the 'coordination-square'. If a plane of best fit is calculated for the *whole* biuret residue, the atom N<sub>2</sub> is found to lie 0.06 Å (about ten times the coordinate standard deviation) away from it. Even with this approximation, however, the complex is bent by 2° 30' along the N<sub>1</sub>N<sub>3</sub> edges of the coordination-square.

In Table 6, we list the coefficients for the planes fitted to various sets of atoms by the least-squares method of Schomaker, Waser, Marsh & Bergman (1959). The table also includes the distances of the atoms from the planes fitted to them, and the angles between the reciprocal axes and the normals to two of the least-squares planes.

#### Atoms not involved in complex

The two independent water molecules and the potassium ion fit into the structure in interesting ways without, however, suggesting that any general conclusions are to be drawn from them.

The potassium ion is in a hole in the structure, its nearest neighbours being the six oxygen atoms shown in the list of K<sup>+</sup>-O distances in Table 5. The six oxygen atoms form a slightly distorted trigonal prism with the potassium ion almost at the centre. Successive prisms in the  $x$ -direction share their triangular faces. The O-K<sup>+</sup>-O angles subtended at the potassium ion by the long edges of the prism all lie between 87.0° and 88.7°, and those angles subtended by the edges of the triangular faces between 75.0° and 79.1°.

One corner of each of these triangular faces is occupied by a water molecule whose oxygen atom is

Table 7. *Parameters of thermal vibration ellipsoids*

$q_i$ , magnitudes of principal axes;  $g_{ik}$ , direction cosines with respect to reciprocal axes  $k$ ;  $\varphi_i$ , angles between principal axes and normal to plane of best fit through biuret molecule (for Cu,  $\varphi_i$  is referred to normal to coordination square)

Atom	Axis $i$	$q_i$	$g_{i1}$	$g_{i2}$	$g_{i3}$	$\varphi_i$
C <sub>1</sub>	1	0.674	0.861	0.135	0.551	63.2°
	2	0.506	0.066	0.927	-0.364	79.8
	3	0.247	-0.504	0.350	0.751	151.1
C <sub>2</sub>	1	0.636	0.630	0.090	0.814	84.1
	2	0.510	-0.521	0.785	0.296	132.9
	3	0.222	0.576	0.613	-0.500	43.6
N <sub>1</sub>	1	1.011	0.768	-0.173	-0.560	8.1
	2	0.492	0.581	0.593	0.598	81.8
	3	0.334	-0.269	0.787	-0.573	89.6
N <sub>2</sub>	1	0.812	0.762	0.039	-0.591	10.4
	2	0.519	0.629	0.196	0.796	83.7
	3	0.372	-0.156	0.980	-0.136	98.3
N <sub>3</sub>	1	0.694	0.977	-0.021	-0.141	20.5
	2	0.512	-0.129	0.732	-0.677	79.1
	3	0.345	0.169	0.681	0.723	107.1
O <sub>1</sub>	1	1.131	0.815	0.217	-0.478	17.2
	2	0.658	-0.388	0.893	-0.255	106.1
	3	0.098	0.431	0.395	0.824	96.0
O <sub>2</sub>	1	1.097	0.941	0.070	-0.263	15.7
	2	0.442	0.093	0.887	0.457	103.7
	3	0.291	0.325	-0.456	0.850	97.6
O <sub>3</sub>	1	1.240	0.987	-0.122	0.179	38.6
	2	0.752	-0.050	0.409	0.906	124.3
	3	0.583	0.156	0.905	-0.385	74.8
O <sub>4</sub>	1	1.202	0.865	0.299	0.465	60.9
	2	0.931	-0.030	0.833	-0.553	78.3
	3	0.610	-0.502	0.466	0.692	148.2
K <sup>+</sup>	1	0.701	0.816	0.325	-0.419	23.7
	2	0.632	0.148	-0.917	-0.359	66.6
	3	0.425	0.559	-0.232	0.834	86.4
Cu	1	0.767	0.896	0.077	-0.372	13.5
	2	0.347	-0.058	0.997	0.527	101.6
	3	0.289	0.440	-0.026	0.927	96.9

of the type O<sub>3</sub>. As a consequence of the alternation of potassium ions and triangles of oxygen atoms, each O<sub>3</sub> has two K<sup>+</sup>-ions as near neighbours at distances of 2.76 and 2.81 Å. An approximate tetrahedron about O<sub>3</sub> is completed by the two atoms, O<sub>4</sub> and O<sub>2</sub> to which it is hydrogen-bonded. [It should be noted by reference to Table 5 that O<sub>3</sub> is not one of the near neighbours of the potassium ion labelled K<sup>+</sup> in Fig. 3, but is so related to the ions labelled K<sup>3'</sup> and K<sup>4'</sup>.]

The oxygen atom O<sub>4</sub> of the second water-molecule also has an approximately tetrahedral environment. It is hydrogen-bonded to O<sub>3</sub>, O<sub>1</sub> and N<sub>3</sub><sup>5</sup>, its hydrogen atoms lying on or near the lines joining it to the latter two.

#### *Analysis of thermal parameters*

The parameters  $q_i$  and  $g_{ik}$  of the thermal vibration ellipsoids for all atoms (other than hydrogens) are shown in Table 7. They were derived from the thermal coefficients listed in Table 2 by the procedure of Rollett & Davies (1955) and represent, respectively, the magnitudes of the principal axes of the vibration ellipsoids and their direction cosines with respect to the reciprocal axes.

We have also calculated the angles between the same principal axes and the normal to the plane of best fit of the biuret residue (plane 4 in Table 6); for the copper atom the direction cosines are referred to the normal to the coordination-square.

The results are in agreement with expectation. The atoms of the complex ion have their maximum amplitudes of vibration roughly normal to the plane of the complex. For the water molecules and potassium ion no such relationship is expected or found. The heavy atoms, copper and potassium, have relatively small amplitudes; the two loosely bound water molecules have the largest amplitudes of vibration. In the biuret residue, the two carbon atoms are exceptional in that, of all the major vibration-axes, theirs are the shortest and do not lie near the normal to the plane of the molecule.

It seems as though, superimposed on the general vibration of the complex as a whole, there is a 'folding' mode in which the copper atom and the two carbon atoms of each biuret residue are stationary and the complex bends along the N<sub>1</sub>N<sub>3</sub>-edges of the coordination-square. In this mode the atoms O<sub>1</sub><sup>'</sup>, N<sub>2</sub><sup>'</sup>, O<sub>2</sub><sup>'</sup>, N<sub>1</sub> and N<sub>3</sub> move 'up' while the atoms N<sub>1</sub><sup>'</sup>, N<sub>3</sub><sup>'</sup>, O<sub>1</sub>, N<sub>2</sub> and O<sub>2</sub> move 'down', and vice versa.

### 5. Colour

To a first approximation, the copper atom in this crystal experiences a tetragonal ligand field. The situation is analogous to that in copper phthalocyanine, for which an energy-level scheme compatible with electron resonance measurements has been proposed by Gibson, Ingram & Schonland (1958). Light absorption in the metal phthalocyanines is, however, a function of the conjugated bond system rather than of the metal atom, so that the analogy cannot be taken further.

The absorption spectrum of crystalline potassium *bis*-biureto cuprate(II) tetrahydrate has a maximum at 505 m $\mu$  and is otherwise almost identical with the spectrum of the compound in aqueous solution between 375 m $\mu$  and 640 m $\mu$  (Dr J. Ferguson, private communication). The complex therefore preserves its structure in solution, so that it is possible to confirm that light absorption at  $\lambda_{\max.} = 505$  m $\mu$  is associated with copper(II) coordinated by four peptide nitrogen atoms (Kober & Haw, 1916; Plekhan, 1952; Kato, 1959).

The amide nitrogen atoms, after the ionization of the protons, are not only negatively charged but are presumably also more polarizable than the nitrogen atoms in ammonia and ethylene diamine. It is therefore not necessary to invoke stabilization by  $\pi$ -bonding (Kato, 1959) to explain why the copper(II) complex of biuret absorbs at a shorter wave-length than do  $\text{Cu}(\text{NH}_3)_4^{++}$  ( $\lambda_{\max.} = 600$  m $\mu$ ) and  $\text{Cu en}_2^+$  ( $\lambda_{\max.} = 550$  m $\mu$ ).

### 6. Discussion

The conclusions to which this structure analysis leads us are:

- (1) In the complex formed by biuret and copper in alkaline solution, the copper atom is bonded to the nitrogen atoms of the amide groups. This confirms the structure first proposed for this complex by Tschugaeff (1907), and supports the view (Dobbie & Kermack, 1955; Datta & Rabin, 1956; Rabin, 1956, 1958) that in alkaline solution metal-peptide binding at peptide nitrogen atoms follows the ionization of the peptide hydrogen atoms.
- (2) The ligand molecule in this complex retains its essential dimensions after coordination. In particular, the dimensions and planarity of the amide-groups are unchanged. The amide-group resonance energy is consequently not reduced by coordination.
- (3) On the other hand, distortion to accommodate the metal atom occurs by the twisting of the amide groups relative to each other, and by the variability of the internal bond angles at the ligand nitrogen atoms.

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### References

- BERGHUIS, J., HAANAPPEL, IJ. M., POTTERS, M., LOOPSTRA, B. O., MACGILLAVRY, C. H. & VEENENDAAL, A. L. (1955). *Acta Cryst.* **8**, 478.
- BRÜCKE, E. (1883). *Monatsh. d. Chemie*, **4**, 203.
- COCHRAN, W. (1951). *Acta Cryst.* **4**, 81, 408.
- CRUICKSHANK, D. W. J. & ROBERTSON, A. P. (1953). *Acta Cryst.* **6**, 698.
- DATTA, S. P. & RABIN, B. R. (1956a). *Trans. Faraday Soc.* **52**, 1123.
- DATTA, S. P. & RABIN, B. R. (1956b). *Biochem. Biophys. Acta*, **19**, 572.
- DAUBEN, C. & TEMPLETON, D. (1955). *Acta Cryst.* **8**, 841.
- DOBBIE, H. & KERMAK, W. O. (1955). *Biochem. J.* **59**, 246, 257.
- FREEMAN, H. C., SMITH, J. E. W. L. & TAYLOR, J. C. (1959). *Nature, Lond.* **184**, 707.
- FREEMAN, H. C. (1958). *Aust. J. Chem.* **11**, 99.
- GIBSON, J. F., INGRAM, D. J. E. & SCHONLAND, D. (1958). *Disc. Faraday Soc.* **26**, 77, 94.
- HOERNI, J. A. & IBERS, J. A. (1954). *Acta Cryst.* **7**, 744.
- HUGHES, E. W. (1941). *J. Amer. Chem. Soc.* **63**, 1737.
- HUGHES, E. W., YAKEL, H. L. & FREEMAN, H. C. (1961). *Acta Cryst.* **14**, 345.
- JOEL, N., VERA, R. & GARAYCOCHEA, I. (1953). *Acta Cryst.* **6**, 1954.
- KATO, M. (1959). *Z. anorg. Chem.* **300**, 84.
- KOBER, P. A. & HAW, A. B. (1916). *J. Amer. Chem. Soc.* **38**, 457.
- LEY, H. & WERNER, F. (1913). *Ber. dtsh. Chem. Ges.* **46**, 4040.
- MATHIESON, A. MCL. (1957). *Acta Cryst.* **10**, 351.
- PLEKHAN, M. I. (1952). *Zhur. Obshchei. Khim.* **22**, 1633.
- RABIN, B. R. (1956). *Trans. Faraday Soc.* **52**, 1134.
- RABIN, B. R. (1958). *Metals and Enzyme Activity*, p. 21. Cambridge: University Press.
- RISING, M. M., HICKS, J. & MOERKE, G. (1930). *J. Biol. Chem.* **89**, 1.
- ROBERTSON, J. M. (1935). *J. Chem. Soc.* p. 615.
- ROBERTSON, J. M. (1936). *J. Chem. Soc.* p. 1195.
- ROBERTSON, J. M. (1937). *J. Chem. Soc.* p. 219.
- ROLLETT, J. S. & DAVIES, D. R. (1955). *Acta Cryst.* **8**, 125.
- SCHIFF, H. (1896). *Ber. dtsh. Chem. Ges.* **29**, 298.
- SCHIFF, H. (1898). *Ann. Chem.* **299**, 238.
- SCHOMAKER, V., WASER, J., MARSH, R. E. & BERGMAN, G. (1959). *Acta Cryst.* **12**, 600.
- TSCHUGAEFF, L. A. (1907). *Ber. dtsh. Chem. Ges.* **40**, 1973.
- WILSON, A. J. C. (1942). *Nature, Lond.* **150**, 152.